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An Existence Result for a Class of Chemically Reacting Systems with Sign-Changing Weights

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Abstract. We prove the existence of positive solutions for the system

$\int -\Delta_p u = \lambda a(x) f(v) u^{-\alpha},$	$x \in \Omega$,
$\left\{-\Delta_{q}v=\lambda b(x)g(u)v^{-\beta}\right\}$	$x \in \Omega$,
u=v=0,	$x \in \partial \Omega$,

where $\Delta_r z = \operatorname{div}(|\nabla z|^{r-2}\nabla z)$, for r > 1 denotes the r-Laplacian operator and λ is a positive parameter, Ω is a bounded domain in \mathbb{R}^n , $n \ge 1$ with sufficiently smooth boundary and $\alpha, \beta \in (0,1)$. Here a(x) and b(x) are C^1 sign-changing functions that maybe negative near the boundary and f,g are C^1 nondecreasing functions, such that $f,g: [0,\infty) \to [0,\infty)$; f(s) > 0, g(s) > 0 for s > 0, $\lim_{s\to\infty} g(s) = \infty$ and

$$\lim_{s\to\infty}\frac{f(Mg(s)^{\frac{1}{q-1}})}{s^{p-1+\alpha}}=0,\qquad\forall M>0.$$

We discuss the existence of positive weak solutions when f, g, a(x) and b(x) satisfy certain additional conditions. We employ the method of sub-supersolution to obtain our results.

AMS Subject Classifications: 35J55, 35J65 Chinese Library Classifications: O175.25, O175.8 Key Words: Positive solutions; chemically reacting systems; sub-supersolutions.

1 Introduction

In this paper, we consider the existence of positive weak solutions for the nonlinear singular system

$$\begin{cases} -\Delta_p u = \lambda a(x) \frac{f(v)}{u^{\alpha}}, & x \in \Omega, \\ -\Delta_q v = \lambda b(x) \frac{g(u)}{v^{\beta}}, & x \in \Omega, \\ u = v = 0, & x \in \partial\Omega, \end{cases}$$
(1.1)

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where $\Delta_r z = \operatorname{div}(|\nabla z|^{r-2}\nabla z)$, for r > 1 denotes the r-Laplacian operator and Ω is a bounded domain in \mathbb{R}^n , $n \ge 1$ with smooth boundary, $\alpha, \beta \in (0,1)$. Here a(x) and b(x) are C^1 sing-changing functions that maybe negative near the boundary and f,g are C^1 nondecreasing functions such that $f,g:[0,\infty) \to [0,\infty)$; f(s) > 0, g(s) > 0 for s > 0.

Systems of singular equations like (1.1) are the stationary counterpart of general evolutionary problems of the form

$$\begin{cases} u_t = \eta \Delta_p u + \lambda \frac{f(v)}{u^{\alpha}}, & x \in \Omega, \\ v_t = \delta \Delta_q v + \lambda \frac{g(u)}{v^{\beta}}, & x \in \Omega, \\ u = v = 0, & x \in \partial \Omega \end{cases}$$

where η and δ are positive parameters. This system is motivated by an interesting applications in chemically reacting systems, where *u* represents the density of an activator chemical substance and *v* is an inhibitor. The slow diffusion of *u* and the fast diffusion of *v* is translated into the fact that η is small and δ is large (see [1]).

Recently, such problems have been studied in [2–4]. Also in [2], the authors have studied the existence results for the system (1.1) in the case $a \equiv 1$, $b \equiv 1$. Here we focus on further extending the study in [2] to the system (1.1). In fact, we study the existence of positive solution to the system (1) with sign-changing weight functions a(x),b(x). Due to these weight functions, the extensions are challenging and nontrivial. Our approach is based on the method of sub-super solutions (see [5–7]).

To precisely state our existence result we consider the eigenvalue problem

$$\begin{cases} -\Delta_r \phi = \lambda |\phi|^{r-2} \phi, & x \in \Omega, \\ \phi = 0, & x \in \partial \Omega. \end{cases}$$
(1.2)

Let $\phi_{1,r}$ be the eigenfunction corresponding to the first eigenvalue $\lambda_{1,r}$ of (1.2) such that $\phi_{1,r}(x) > 0$ in Ω , and $\|\phi_{1,r}\|_{\infty} = 1$ for r = p,q. Let $m,\sigma,\delta > 0$ be such that

$$\sigma \le \phi_{1,r} \le 1, \quad x \in \Omega - \overline{\Omega_{\delta}}, \tag{1.3}$$

$$\lambda_{1,r}\phi_{1,r}^{r} - \left(1 - \frac{sr}{r - 1 + s}\right) |\nabla\phi_{1,r}|^{r} \le -m, \quad x \in \overline{\Omega_{\delta}}, \tag{1.4}$$

for r=p,q, and $s=\alpha,\beta$, where $\overline{\Omega_{\delta}}:=\{x\in\Omega|d(x,\partial\Omega)\leq\delta\}$. (This is possible since $|\nabla\phi_{1,r}|^r\neq 0$ on $\partial\Omega$ while $\phi_{1,r}=0$ on $\partial\Omega$ for r=p,q. We will also consider the unique solution $\zeta_r\in W_0^{1,r}(\Omega)$ of the boundary value problem

$$\begin{cases} -\Delta_r \zeta_r = 1, & x \in \Omega, \\ \zeta_r = 0, & x \in \partial \Omega \end{cases}$$

to discuss our existence result, it is known that $\zeta_r > 0$ in Ω and $\partial \zeta_r / \partial n < 0$ on $\partial \Omega$.

Here we assume that the weight functions a(x) and b(x) take negative values in $\overline{\Omega_{\delta}}$, but require a(x) and b(x) be strictly positive in $\Omega - \overline{\Omega_{\delta}}$. To be precise we assume that there exist positive constants a_0 , a_1 , b_0 and b_1 such that $a(x) \ge -a_0$, $b(x) \ge -b_0$ on $\overline{\Omega_{\delta}}$ and $a(x) \ge a_1$, $b(x) \ge b_1$ on $\Omega - \overline{\Omega_{\delta}}$.

2 Existence result

In this section, we shall establish our existence result via the method of sub - supersolution. A pair of nonnegative functions $(\psi_1, \psi_2) \in W^{1,p} \cap C(\overline{\Omega}) \times W^{1,q} \cap C(\overline{\Omega})$ and $(z_1, z_2) \in W^{1,p} \cap C(\overline{\Omega}) \times W^{1,q} \cap C(\overline{\Omega})$ are called a subsolution and supersulution of (1.1) if they satisfy $(\psi_1, \psi_2) = (0, 0) = (z_1, z_2)$ on $\partial\Omega$ and

$$\begin{split} &\int_{\Omega} |\nabla \psi_{1}|^{p-2} \nabla \psi_{1} \cdot \nabla w \, \mathrm{d}x \leq \lambda \int_{\Omega} a(x) \frac{f(\psi_{2})}{\psi_{1}^{\alpha}} w \, \mathrm{d}x, \\ &\int_{\Omega} |\nabla \psi_{2}|^{q-2} \nabla \psi_{2} \cdot \nabla w \, \mathrm{d}x \leq \lambda \int_{\Omega} b(x) \frac{g(\psi_{1})}{\psi_{2}^{\beta}} w \, \mathrm{d}x, \\ &\int_{\Omega} |\nabla z_{1}|^{p-2} \nabla z_{1} \cdot \nabla w \, \mathrm{d}x \geq \lambda \int_{\Omega} a(x) \frac{f(z_{2})}{z_{1}^{\alpha}} w \, \mathrm{d}x, \\ &\int_{\Omega} |\nabla z_{2}|^{q-2} \nabla z_{2} \cdot \nabla w \, \mathrm{d}x \geq \lambda \int_{\Omega} b(x) \frac{g(z_{1})}{z_{2}^{\beta}} w \, \mathrm{d}x, \end{split}$$

for all $w \in W = \{w \in C_0^{\infty}(\Omega) | w \ge 0, x \in \Omega\}$. Then the following result holds:

Lemma 2.1. (See [5]) Suppose there exist sub and supersolutions (ψ_1, ψ_2) and (z_1, z_2) respectively of (1.1) such that $(\psi_1, \psi_2) \leq (z_1, z_2)$. Then (1.1) has a solution (u, v) such that $(u, v) \in [(\psi_1, \psi_2), (z_1, z_2)]$.

To state our results precisely we introduce the following hypotheses :

- **(H1)** $f,g:[0,\infty) \to [0,\infty)$ are C^1 nondecreasing functions such that f(s), g(s) > 0 for s > 0, and $\lim_{s\to\infty} g(s) = \infty$.
- (H2) $\lim_{s\to\infty} \frac{f(Mg(s)^{\frac{1}{q-1}})}{s^{p-1+\alpha}} = 0$, for all M > 0.
- (H3) Suppose that there exists $\epsilon > 0$ such that:

$$\frac{e^{\frac{p-1+\alpha}{p-1}}\lambda_{1,p}}{a_1f(\frac{q-1+\beta}{q}e^{\frac{1}{q-1}}\sigma^{\frac{q}{q-1+\beta}})} \leq \min\left\{\frac{me^{\frac{\alpha+p-1}{p-1}}(\frac{p-1+\alpha}{p})^{\alpha}}{a_0f(e^{\frac{1}{q-1}})}, \frac{me^{\frac{\beta+q-1}{q-1}}(\frac{q-1+\beta}{q})^{\beta}}{b_0g(e^{\frac{1}{p-1}})}\right\}, \\
\frac{e^{\frac{q-1+\beta}{q-1}}\lambda_{1,q}}{b_1g(\frac{p-1+\alpha}{p}e^{\frac{1}{p-1}}\sigma^{\frac{p}{p-1+\alpha}})} \leq \min\left\{\frac{me^{\frac{\alpha+p-1}{p-1}}(\frac{p-1+\alpha}{p})^{\alpha}}{a_0f(e^{\frac{1}{q-1}})}, \frac{me^{\frac{\beta+q-1}{q-1}}(\frac{q-1+\beta}{q})^{\beta}}{b_0g(e^{\frac{1}{p-1}})}\right\}.$$

We are now ready to give our existence result .

Theorem 2.1. Assume that

- (a) $p \ge n$ or (b) p < n and $\alpha < \frac{p}{n}$,
- (c) $q \ge n$ or (d) q < n and $\beta < \frac{q}{n}$.

Let (H1)-(H3) *hold. Then there exists a positive solution of* (1.1) *for every* $\lambda \in [\lambda_*(\epsilon), \lambda^*(\epsilon)]$ *, where*

$$\lambda^* = \min\left\{\frac{m\epsilon^{\frac{\alpha+p-1}{p-1}}(\frac{p-1+\alpha}{p})^{\alpha}}{a_0f(\epsilon^{\frac{1}{q-1}})}, \frac{m\epsilon^{\frac{\beta+q-1}{q-1}}(\frac{q-1+\beta}{q})^{\beta}}{b_0g(\epsilon^{\frac{1}{p-1}})}\right\},$$

and

$$\lambda_* = \max\left\{\frac{\epsilon^{\frac{\alpha+p-1}{p-1}}\lambda_{1,p}}{a_1f(\frac{q-1+\beta}{q}\epsilon^{\frac{1}{q-1}}\sigma^{\frac{q}{q-1+\beta}})}, \frac{\epsilon^{\frac{\beta+q-1}{q-1}}\lambda_{1,q}}{b_1g(\frac{p-1+\beta}{p}\epsilon^{\frac{1}{p-1}}\sigma^{\frac{p}{p-1+\alpha}})}\right\}.$$

Remark 2.1. Note that (*H*3) implies $\lambda_* < \lambda^*$.

Example 2.1. Let $f(s) = s^4 - 1$, g(s) = s - 1 and $(p = q = 3, n < 6, \alpha = 1/2, \beta = 1/2)$. Here f(s), g(s) > 0 for s > 0, f, g are non-decreasing functions and

$$\lim_{s \to \infty} \frac{f(Mg(s)^{\frac{1}{2}})}{s^{\frac{5}{2}}} = 0,$$

for all M > 0, and $\lim_{s\to\infty} g(s) = \infty$. We can choose $\epsilon > 0$ so small that f, g satisfy (H3).

Proof of Theorem 2.1 We shall verify that

$$(\psi_1,\psi_2) = \left(\frac{p-1+\alpha}{p} e^{\frac{1}{p-1}} \phi_{1,p}^{\frac{p}{p-1+\alpha}}, \quad \frac{q-1+\beta}{q} e^{\frac{1}{q-1}} \phi_{1,q}^{\frac{q}{q-1+\beta}}\right),$$

is a sub-solution of (1.1). Let $w \in W$. Then a calculation shows that

$$\nabla\psi_1 = \epsilon^{\frac{1}{p-1}} \nabla\phi_{1,p} \phi_{1,p}^{\frac{1-\alpha}{p-1+\alpha}},$$

and we have

$$\begin{split} \int_{\Omega} |\nabla \psi_{1}|^{p-2} \nabla \psi_{1} \cdot \nabla w \, \mathrm{d}x = & \epsilon \int_{\Omega} \phi_{1,p}^{1-\frac{\alpha p}{p-1+\alpha}} |\nabla \phi_{1,p}|^{p-2} \nabla \phi_{1,p} \nabla w \, \mathrm{d}x \\ = & \epsilon \int_{\Omega} |\nabla \phi_{1,p}|^{p-2} \nabla \phi_{1,p} \left\{ \nabla (\phi_{1,p}^{1-\frac{\alpha p}{p-1+\alpha}} w) - w \nabla (\phi_{1,p}^{1-\frac{\alpha p}{p-1+\alpha}}) \right\} \, \mathrm{d}x \\ = & \epsilon \left\{ \int_{\Omega} \left[\lambda_{1,p} \phi_{1,p}^{p-\frac{\alpha p}{p-1+\alpha}} - |\nabla \phi_{1,p}|^{p-2} \nabla \phi_{1,p} \nabla (\phi_{1,p}^{1-\frac{\alpha p}{p-1+\alpha}}) \right] w \, \mathrm{d}x \right\} \end{split}$$

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$$= \epsilon \left\{ \int_{\Omega} \left[\lambda_{1,p} \phi_{1,p}^{p-\frac{\alpha p}{p-1+\alpha}} - |\nabla \phi_{1,p}|^{p} (1-\frac{\alpha p}{p-1+\alpha}) \phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}} \right] w dx \right\}$$
$$= \epsilon \phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}} \left\{ \int_{\Omega} \left[\lambda_{1,p} \phi_{1,p}^{p} - |\nabla \phi_{1,p}|^{p} (1-\frac{\alpha p}{p-1+\alpha}) \right] w dx \right\}.$$

Similarly

$$\int_{\Omega} |\nabla \psi_2|^{q-2} \nabla \psi_2 \cdot \nabla w \, \mathrm{d}x = \epsilon \phi_{1,q}^{-\frac{\beta q}{q-1+\beta}} \left\{ \int_{\Omega} [\lambda_{1,q} \phi_{1,q}^q - |\nabla \phi_{1,q}|^q (1 - \frac{\beta q}{q-1+\beta})] w \, \mathrm{d}x \right\}.$$

First we consider the case when $x \in \overline{\Omega_{\delta}}$. We have

$$\lambda_{1,p}\phi_{1,p}^{p} - (1 - \frac{\alpha p}{p - 1 + \alpha})|\nabla \phi_{1,p}|^{p} \leq -m.$$

Since $\lambda \leq \lambda^*$ then

$$\lambda \leq \frac{m \epsilon^{\frac{\alpha+p-1}{p-1}} (\frac{p-1+\alpha}{p})^{\alpha}}{a_0 f(\epsilon^{\frac{1}{q-1}})}.$$

Hence

$$\begin{split} & \varepsilon \phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}} \big(\lambda_{1,p} \phi_{1,p}^{p} - (1 - \frac{\alpha p}{p-1+\alpha}) |\nabla \phi_{1,p}|^{p} \big) \leq -m \varepsilon \phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}} \\ & \leq -\lambda a_{0} \frac{f(\varepsilon^{\frac{1}{q-1}}) (\frac{p-1+\alpha}{p})^{-\alpha} \phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}}}{\varepsilon^{\frac{\alpha}{p-1}}} \leq -\lambda a_{0} \frac{f(\frac{q-1+\beta}{q} \varepsilon^{\frac{1}{q-1}} \phi_{1,q}^{\frac{1}{q-1}+\beta}) (\frac{p-1+\alpha}{p})^{-\alpha} \phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}}}{\varepsilon^{\frac{\alpha}{p-1}}} \\ & \leq \lambda a(x) \frac{f(\psi_{2})}{\psi_{1}^{\alpha}}. \end{split}$$

A similar argument shows that

$$-m\epsilon\phi_{1,q}^{-\frac{\beta q}{q-1+\beta}}\left(\lambda_{1,q}\phi_{1,q}^{q}-\left(1-\frac{\beta q}{q-1+\beta}\right)|\nabla\phi_{1,q}|^{q}\right)\leq\lambda b(x)\frac{g(\psi_{1})}{\psi_{2}^{\beta}}.$$

On the other hand, on $\Omega - \overline{\Omega_{\delta}}$, we have $1 \ge \phi_{1,r} \ge \sigma$ for r = p,q. Also $a(x) \ge a_1$, $b(x) \ge b_1$ and since $\lambda \ge \lambda_*$, we have

$$\lambda \ge \frac{e^{\frac{q+p-1}{p-1}}\lambda_{1,p}}{a_1f(\frac{q-1+\beta}{q}e^{\frac{1}{q-1}}\sigma^{\frac{q}{q-1+\beta}})}.$$

Hence

$$\epsilon \phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}} \left(\lambda_{1,p} \phi_{1,p}^p - \left(1 - \frac{\alpha p}{p-1+\alpha}\right) |\nabla \phi_{1,p}|^p \right) \leq \epsilon \lambda_{1,p} \phi_{1,p}^{p-\frac{\alpha p}{p-1+\alpha}}$$

$$\leq \epsilon \lambda_{1,p} \phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}} \leq \lambda a_1 \frac{f(\frac{q-1+\beta}{q}\epsilon^{\frac{1}{q-1}}\sigma^{\frac{q}{q-1+\beta}})\phi_{1,p}^{-\frac{\alpha p}{p-1+\alpha}}(\frac{p-1+\alpha}{p})^{-\alpha}}{\epsilon^{\frac{\alpha}{p-1}}} \leq \lambda a(x) \frac{f(\psi_2)}{\psi_1^{\alpha}}.$$

A similar argument shows that

$$\epsilon \phi_{1,q}^{-\frac{\beta q}{q-1+\beta}} \left(\lambda_{1,q} \phi_{1,q}^{q} - (1 - \frac{\beta q}{q-1+\beta}) |\nabla \phi_{1,q}|^{q} \right) \leq \lambda b(x) \frac{g(\psi_{1})}{\psi_{2}^{\beta}}$$

Hence

$$\int_{\Omega} |\nabla \psi_1|^{p-2} \nabla \psi_1 \cdot \nabla w \, \mathrm{d}x \leq \lambda \int_{\Omega} a(x) \frac{f(\psi_2)}{\psi_1^{\alpha}} w \, \mathrm{d}x,$$

and

$$\int_{\Omega} |\nabla \psi_2|^{q-2} \nabla \psi_2 \cdot \nabla w \, \mathrm{d}x \leq \lambda \int_{\Omega} b(x) \frac{g(\psi_1)}{\psi_2^{\beta}} w \, \mathrm{d}x.$$

Thus, $(\psi_1, \psi_2) = (\frac{p-1+\alpha}{p} \epsilon^{\frac{1}{p-1}} \phi_{1,p}^{\frac{p}{p-1+\alpha}}, \frac{q-1+\beta}{q} \epsilon^{\frac{1}{q-1}} \phi_{1,q}^{\frac{q}{q-1+\beta}})$ is a positive subsolution of (1.1). Now, we construct a suupersolution $(z_1, z_2) \ge (\psi_1, \psi_2)$. When

(a) $p \ge n$ or (b) p < n and $\alpha < \frac{p}{n}$, (c) $q \ge n$ or (d) q < n and $\beta < \frac{q}{n}$,

from [8], we know that are functions $w_1 \in W_0^{1,p}(\Omega) \cap C(\overline{\Omega})$ and $w_2 \in W_0^{1,q}(\Omega) \cap C(\overline{\Omega})$ such that

$$\begin{cases} -\Delta_p w_1 = \frac{1}{w_1^{\alpha}}, & x \in \Omega, \\ w_1 = 0, & x \in \partial \Omega \end{cases}$$

and

$$\begin{cases} -\Delta_q w_2 = \frac{1}{w_2^{\beta}}, & x \in \Omega, \\ w_2 = 0, & x \in \partial\Omega, \end{cases}$$

are satisfying $w_1 \ge \theta \zeta_p$ and $w_2 \ge \theta \zeta_q$ for some $\theta > 0$. Now, we will prove there exists $c \gg 1$ such that

$$(z_1,z_2) = (cw_1,g(c||w_1||_{\infty})^{\frac{1}{q-1}}w_2),$$

is a supersolution of (1.1). A calculation shows that :

$$\int_{\Omega} |\nabla z_1|^{p-2} \nabla z_1 \cdot \nabla w \, \mathrm{d}x = c^{p-1} \int_{\Omega} |\nabla w_1|^{p-2} \nabla w_1 \cdot \nabla w \, \mathrm{d}x = c^{p-1} \int_{\Omega} \frac{w}{w_1^{\alpha}} \, \mathrm{d}x,$$

by (H2) we know that, for $c \gg 1$,

$$\frac{1}{\lambda \|a(x)\|_{\infty}} \ge \frac{f(\|w_2\|_{\infty}(g(c\|w_1\|_{\infty}))^{\frac{1}{q-1}})}{c^{p-1+\alpha}}.$$

Hence

$$\begin{aligned} \frac{c^{p-1}}{w_1^{\alpha}} \ge &\lambda \|a(x)\|_{\infty} \frac{f(\|w_2\|_{\infty}(g(c\|w_1\|_{\infty}))^{\frac{1}{q-1}})}{(cw_1)^{\alpha}} \\ \ge &\lambda a(x) \frac{f(w_2(g(c\|w_1\|_{\infty}))^{\frac{1}{q-1}})}{(cw_1)^{\alpha}} = &\lambda a(x) \frac{f(z_2)}{z_1^{\alpha}}, \end{aligned}$$

now from (H1), we know that $g(s) \rightarrow \infty$ as $s \rightarrow \infty$. Thus, for $c \gg 1$

$$\frac{\lambda \|b(x)\|_{\infty}}{g(c\|w_1\|_{\infty})^{\frac{\beta}{q-1}}} \leq 1,$$

and we have for $c \gg 1$,

$$\int_{\Omega} |\nabla z_2|^{q-2} \nabla z_2 \cdot \nabla w \, \mathrm{d}x = \int_{\Omega} \frac{g(c \| w_1 \|_{\infty})}{w_2^{\beta}} w \, \mathrm{d}x$$
$$\geq \lambda \| b(x) \|_{\infty} \int_{\Omega} \frac{g(cw_1)}{g(c \| w_1 \|_{\infty})^{\frac{\beta}{q-1}} w_2^{\beta}} w \, \mathrm{d}x \geq \lambda \int_{\Omega} b(x) \frac{g(z_1)}{z_2^{\beta}} w \, \mathrm{d}x,$$

i.e., (z_1, z_2) is a supersolution of (1.1). Furthermore, *c* can be chosen large enough so that $(z_1, z_2) \ge (\psi_1, \psi_2)$, since $g(s) \to \infty$ as $s \to \infty$. Thus, there exist a positive solution (u, v) of (1) such that $(\psi_1, \psi_2) \le (u, v) \le (z_1, z_2)$. This completes the proof of Theorem 2.1.

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