# A NOTE ON THE NONCONFORMING FINITE ELEMENTS FOR ELLIPTIC PROBLEMS* 

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#### Abstract

In this paper, a class of rectangular finite elements for $2 m$-th-oder elliptic boundary value problems in $n$-dimension ( $m, n \geq 1$ ) is proposed in a canonical fashion, which includes the $(2 m-1)$-th Hermite interpolation element $(n=1)$, the $n$-linear finite element ( $m=1$ ) and the Adini element $(m=2)$. A nonconforming triangular finite element for the plate bending problem, with convergent order $\mathcal{O}\left(h^{2}\right)$, is also proposed.


Mathematics subject classification: 65N30.
Key words: Nonconforming finite element, Elliptic boundary value problem, Plate bending problem.

## 1. Introduction

When the conforming finite element is used for numerically discretizing the elliptic problem, the convergence of the numerical solution to the exact solution depends on the approximation of the finite element space only. But the strong continuity requirement makes it difficult to construct such a conforming finite element. The idea of nonconforming finite element lies in that such difficulty can be overcome by loosing the request on the continuity. However, the loss of continuity will bring in the so-called consistent error, and some fundamental continuity of the finite element space is still necessary for well-posedness and convergence. This is the reason that most of the finite elements, conforming or nonconforming, were constructed case by case, depending on the order of the problem and sometimes the dimensions (cf. [1-3, 5, 7, $8,12,14,15,17])$. A unified approach of constructing finite elements for general problems is still of theoretical and practical interest. Recently, a class of finite elements was discussed in a canonical fashion in [16], for all $n$-dimensional $2 m$-th-order elliptic problem with $n \geq m \geq 1$. The well-known nonconforming linear element for the second-order problem and the Morley element for fourth-order problem are examples of this class. The class of finite elements is established on simplices, and makes use of the piecewise polynomials of the lowest degree. The nodal parameters are the natural ones to guarantee the fundamental continuity, and the consistency error can be controlled simultaneously.

[^0]In this paper, we will discuss the choice of nodal parameters that can be used to construct nonconforming finite elements, with admissible consistency error. We will first propose a class of rectangular finite elements for $n$-dimensional $2 m$-th-oder problems $(m, n \geq 1$ ) in a canonical fashion. The degrees of freedom are the values of function and all derivatives up to $(m-1)$ -th-oder at all vertices of $n$-rectangle. The basic fundamental continuity is guaranteed and an $\mathcal{O}(h)$ convergence rate is shown. The $(2 m-1)$-th Hermite interpolation element $(n=1)$, the $n$-linear finite element $(m=1)$ and the Adini element $(m=2)$ all belong to this class.

As almost all of the nonconforming finite elements are convergent in energy norm with order $\mathcal{O}(h)$, and the consistency error is the main limit, we will discuss the possibility of improving the convergence rate by strengthening the continuity of the finite element space. We choose the plate bending problem as an example. There have been successful attempts via other approaches, like conforming finite element, quasi-conforming finite elements (cf. $[4,6,11]$ ) and the double set parameter element (cf. [9]). But most nonconforming element for the plate bending problem, such as the Morley element [8], two Veubake elements [12], the NZT element [14], the rectangle Morley element (cf. [15]) and the Adini element (cf. [1]), are convergent with order $\mathcal{O}(h)$. In this work, a new nonconforming plate element will be given, with a convergence rate of $\mathcal{O}\left(h^{2}\right)$ in energy norm.

Finally, based on the new plate element, a new Zienkiewicz-type element will be deduced and reported for comparison. The new Zienkiewicz-type element is convergent for the plate bending problem with order $\mathcal{O}(h)$. Its consistent error is of order $\mathcal{O}\left(h^{2}\right)$ which is better than the two dimensional Zienkiewicz-type element proposed in [14]. In fact, the phenomenon that the consistency error can perform better than the approximation error has seldom been reported in literatures.

The paper is organized as follows. The rest of this section gives some basic notations. Section 2 gives the description of the class of rectangular finite elements. Section 3 gives the description of the new plate elements. Section 4 shows their convergence. Section 5 gives some numerical results for the new plate element.

Let $n$ be a positive integer. Given a nonnegative integer $k$ and a bounded domain $G \subset R^{n}$ with boundary $\partial G$, let $H^{k}(G), H_{0}^{k}(G),\|\cdot\|_{k, G}$ and $|\cdot|_{k, G}$ denote the usual Sobolev spaces, norm and semi-norm respectively. Let $(\cdot, \cdot)$ denote the inner product of $L^{2}(\Omega)$.

We will use $\alpha, \beta, \gamma$ to denote $n$ dimensional multi-indexes. Define

$$
\partial^{\alpha}=\frac{\partial^{|\alpha|}}{\partial x_{1}^{\alpha_{1}} \cdots \partial x_{n}^{\alpha_{n}}}, \quad|\alpha|=\alpha_{1}+\cdots+\alpha_{n}
$$

A finite element can be represented by a triple $\left(T, P_{T}, D_{T}\right)$ with $T$ the geometric shape, $P_{T}$ the shape function space and $D_{T}$ the vector of degrees of freedom, provided that $D_{T}$ is $P_{T^{-} \text {-unisolvent (see [5]). }}$

Let $\Omega$ be a bounded polyhedron domain of $R^{n}$. For mesh size $h$ with $h \rightarrow 0$, let $\mathcal{T}_{h}$ be a partition of $\Omega$ corresponding to a finite element $\left(T, P_{T}, D_{T}\right)$, and let $V_{h}, V_{h 0}$ be the finite element spaces corresponding to the element and $\mathcal{T}_{h}$. Throughout this paper, we assume that $\left\{\mathcal{T}_{h}\right\}$ is shape regular.

For a subset $B \subset R^{n}$ and a nonnegative integer $r$, let $P_{r}(B)$ be the space of all polynomials defined on $B$ with degree not greater than $r$, and $Q_{r}(B)$ the space of all polynomials with degree in each variable not greater than $r$.

## 2. A Class of Rectangular Finite Elements

Let $m$ be a positive integer. This section is devoted to the rectangular finite element for the $2 m$-th-oder elliptic boundary value problem in $n$-dimension.

Let $T$ be an $n$-rectangle with each edge parallel to some coordinate axis respectively. Then there exist $n$ positive numbers $h_{1}, h_{2}, \cdots, h_{n}$, such that,

$$
\begin{equation*}
T=\left\{x=\left(x_{1}, x_{2}, \cdots, x_{n}\right)^{\mathrm{T}} \mid x_{i}=x_{i}^{0}+\xi_{i} h_{i},-1 \leq \xi_{i} \leq 1,1 \leq i \leq n\right\} \tag{2.1}
\end{equation*}
$$

where $x^{0}$ is the center point of $T$. Define

$$
\begin{equation*}
\xi_{i}=\frac{1}{h_{i}}\left(x_{i}-x_{i}^{0}\right), \quad 1 \leq i \leq n \tag{2.2}
\end{equation*}
$$

and set $\xi=\left(\xi_{1}, \xi_{2}, \cdots, \xi_{n}\right)^{\mathrm{T}}$. Denote $2^{n}$ vertices of $T$ by $a_{j}, 1 \leq j \leq 2^{n}$, and $\left(\xi_{1}, \xi_{2}, \cdots, \xi_{n}\right)^{\mathrm{T}}$ corresponding to $a_{j}$ by $\Xi_{j}=\left(\xi_{1 j}, \xi_{2 j}, \cdots, \xi_{n j}\right)^{\mathrm{T}}$.

For $\xi=\left(\xi_{1}, \xi_{2}, \cdots, \xi_{n}\right)^{\mathrm{T}}$ and multi-index $\alpha=\left(\alpha_{1}, \alpha_{2}, \cdots, \alpha_{n}\right)$, define

$$
\xi^{\alpha}=\prod_{i=1}^{n} \xi_{i}^{\alpha_{i}}, \quad \alpha!=\prod_{i=1}^{n} \alpha_{i}!.
$$

Set

$$
\begin{align*}
& \varphi_{j}=\frac{1}{2^{n}} \prod_{i=1}^{n}\left(1+\xi_{i j} \xi_{i}\right), \quad 1 \leq j \leq 2^{n}  \tag{2.3}\\
& P_{T, m}=\operatorname{span}\left\{\varphi_{j} \xi^{2 \alpha}\left|1 \leq j \leq 2^{n},|\alpha|<m\right\}\right. \tag{2.4}
\end{align*}
$$

Then $\varphi_{j}, 1 \leq j \leq 2^{n}$, form a basis of $Q_{1}(T)$, and

$$
\begin{equation*}
P_{T, m}=\operatorname{span}\left\{p \xi^{2 \alpha}\left|p \in Q_{1}(T),|\alpha|<m\right\}\right. \tag{2.5}
\end{equation*}
$$

The rectangular finite element of order $m$ is defined by the triple $\left(T, P_{T}, D_{T}\right)$ as follows,

1. $T$ is the $n$-rectangle described by (2.1);
2. $P_{T}=P_{T, m}$;
3. the components of $D_{T}(v)$ for any $v \in C^{m-1}(T)$ are

$$
\partial^{\alpha} v\left(a_{j}\right), \quad|\alpha|<m, \quad 1 \leq j \leq 2^{n}
$$

Lemma 2.1. For the rectangular finite element of order $m, D_{T}$ is $P_{T}$-unisolvent and $P_{2 m-1}(T)$ $\subset P_{T}$.

Proof. It is obvious that the dimensions of $D_{T}$ and $P_{T}$ are all $2^{n} C_{n+m-1}^{m-1}$. For $1 \leq j \leq 2^{n}$, set

$$
\begin{equation*}
\varphi_{j, \alpha}=\frac{\Xi_{j}^{\alpha}}{\alpha!2^{|\alpha|}}\left(\xi_{1}^{2}-1\right)^{\alpha_{1}}\left(\xi_{2}^{2}-1\right)^{\alpha_{2}} \cdots\left(\xi_{n}^{2}-1\right)^{\alpha_{n}} \varphi_{j}, \quad|\alpha|<m \tag{2.6}
\end{equation*}
$$

Write the partial derivative with respect to $\xi$ as

$$
\partial_{\xi}^{\alpha}=\frac{\partial^{|\alpha|}}{\partial \xi_{1}^{\alpha_{1}} \cdots \partial \xi_{n}^{\alpha_{n}}}
$$

It can be verified that

$$
\partial_{\xi}^{\beta} \varphi_{j, \alpha}\left(a_{k}\right)=\left\{\begin{array}{ll}
1, & \beta=\alpha \text { and } j=k,  \tag{2.7}\\
0, & \text { otherwise },
\end{array} \quad 1 \leq j, \quad k \leq 2^{n}, \quad|\beta| \leq|\alpha|<m .\right.
$$

Define

$$
\begin{equation*}
\psi_{j, \alpha}=\varphi_{j, \alpha} \tag{2.8}
\end{equation*}
$$

when $|\alpha|=m-1$, and

$$
\begin{equation*}
\psi_{j, \alpha}=\varphi_{j, \alpha}-\sum_{k=1}^{2^{n}} \sum_{|\beta|=|\alpha|+1}^{m-1} \partial_{\xi}^{\beta} \varphi_{j, \alpha}\left(a_{k}\right) \psi_{k, \beta} \tag{2.9}
\end{equation*}
$$

when $|\alpha|<m-1$. Then

$$
\partial_{\xi}^{\beta} \psi_{j, \alpha}\left(a_{k}\right)=\left\{\begin{array}{ll}
1, & \beta=\alpha \text { and } j=k,  \tag{2.10}\\
0, & \text { otherwise },
\end{array} \quad 1 \leq j, k \leq 2^{n},|\alpha|,|\beta|<m .\right.
$$

Therefore, $h_{1}^{\alpha_{1}} \cdots h_{n}^{\alpha_{n}} \psi_{j, \alpha}\left(1 \leq j \leq 2^{n},|\alpha|<m\right)$ are the basis functions corresponding to the degrees of freedom since $\partial_{\xi}^{\alpha}=h_{1}^{\alpha_{1}} \cdots h_{n}^{\alpha_{n}} \partial^{\alpha}$. Thus, we obtain that $D_{T}$ is $P_{T}$-unisolvent.

Now we show that $P_{2 m-1}(T) \subset P_{T, m}$. Let $p \in P_{2 m-1}(T)$, then $p$ can be written as

$$
p=\sum_{|\alpha| \leq 2 m-1} C_{\alpha} x^{\alpha},
$$

with $C_{\alpha}$ constants. For term $C_{\alpha} x^{\alpha}$ with $|\alpha| \leq 2 m-1$, define $\beta$ and $\gamma$ by

$$
\beta_{i}=\left\{\begin{array}{ll}
\frac{\alpha_{i}}{2}, & \alpha_{i} \text { is even, } \\
\frac{\alpha_{i}-1}{2}, & \alpha_{i} \text { is odd, }
\end{array} \quad \gamma_{i}=\left\{\begin{array}{ll}
0, & \alpha_{i} \text { is even, } \\
1, & \alpha_{i} \text { is odd, }
\end{array} \quad 1 \leq i \leq n\right.\right.
$$

Then $\alpha=2 \beta+\gamma$, so that $C_{\alpha} x^{\alpha}=C_{\alpha} x^{\gamma} x^{2 \beta} \in P_{T, m}$ by (2.5) and the fact that $x^{\gamma} \in Q_{1}(T)$ and $|\beta|<m$. Consequently, $p \in P_{T, m}$.

For the rectangular finite element of order $m$, the corresponding finite element spaces $V_{h}$ and $V_{h 0}$ are defined as follows. $V_{h}=\left\{v \in L^{2}(\Omega)|v|_{T} \in P_{T, m}, \forall T \in T_{h}, \partial^{\alpha} v,|\alpha|<m\right.$, are continuous at all vertices of elements in $\left.\mathcal{T}_{h}\right\} . V_{h 0}=\left\{v \in V_{h}\left|\partial^{\alpha} v,|\alpha|<m\right.\right.$, vanish at all vertices of elements in $\mathcal{T}_{h}$ which are belonging to $\left.\partial \Omega\right\}$.

Remark 2.1. The rectangular finite element of order $m$ is just the $(2 m-1)$-th-oder Hermite interpolation element when $n=1$, the $n$-linear finite element when $m=1$ and the Adini element when $m=2$. For the $2 m$-th-oder problems, the rectangular finite element of order $m$ is conforming when $m=1$ or $n=1$, otherwise it is nonconforming. The rectangular finite element of order $m$ can be viewed as the natural and reasonable generalizations of the one dimensional $(2 m-1)$-th-oder Hermite interpolation element to higher dimensions or the $n$ linear finite element to higher order problems. This generalization shows that the conforming elements and the nonconforming elements are in same category.

Now let $\Pi_{T, m}$ be the corresponding interpolation operator to the rectangular finite element of order $m$.

Lemma 2.2. For the rectangular finite element of order m,

$$
\begin{equation*}
\int_{T} \frac{\partial}{\partial x_{i}}\left(\partial^{\beta} p-\Pi_{T, 1} \partial^{\beta} p\right) \mathrm{d} x=0, \quad 1 \leq i \leq n, \quad|\beta|<m, \quad \forall p \in P_{T, m} . \tag{2.11}
\end{equation*}
$$

Proof. Let $p \in P_{T, m}, 1 \leq i \leq n$ and $|\beta|<m$. We know by (2.4) that $\partial^{\beta} p$ is a linear combination of the following functions,

$$
F_{j, \alpha}=\partial^{\beta}\left(\left(\xi_{1}^{2}-1\right)^{\alpha_{1}}\left(\xi_{2}^{2}-1\right)^{\alpha_{2}} \cdots\left(\xi_{n}^{2}-1\right)^{\alpha_{n}} \varphi_{j}\right), \quad 1 \leq j \leq 2^{n}, \quad|\alpha|<m
$$

For $1 \leq j \leq 2^{n}$ and $|\alpha|<m$, set

$$
f_{i}=\frac{\mathrm{d}^{\beta_{i}}}{\mathrm{~d} \xi_{i}^{\beta_{i}}}\left(\xi_{i}^{2}-1\right)^{\alpha_{i}}, \quad g_{i}=\frac{\mathrm{d}^{\beta_{i}}}{\mathrm{~d} \xi_{i}^{\beta_{i}}}\left(\xi_{i}\left(\xi_{i}^{2}-1\right)^{\alpha_{i}}\right) .
$$

Then $F_{j, \alpha}$ can be written as the sum of such terms that each term has two factors, one is $f_{i}$ or $g_{i}$, and another is independent of component $\xi_{i}$. Define

$$
G_{i}(T)=\left\{v \in C^{\infty}(T) \left\lvert\, \int_{T} \frac{\partial v}{\partial \xi_{i}} \mathrm{~d} x=0\right., v\left(a_{j}\right)=0,1 \leq j \leq 2^{n}\right\}
$$

a) $\beta_{i}<\alpha_{i}$. In this case, $\frac{\mathrm{d} f_{i}}{\mathrm{~d} \xi_{i}}$ is just the Legendre polynomial of $\xi_{i}$ or its integral. Hence

$$
\int_{-1}^{1} \frac{\mathrm{~d} f_{i}}{\mathrm{~d} \xi_{i}} \mathrm{~d} \xi_{i}=0
$$

On the other hand, $f_{i}$ vanishes when $\xi_{i}= \pm 1$. Then $f_{i} \in G_{i}(T)$.
b) $\beta_{i} \geq \alpha_{i}$ and $\alpha_{i} \leq 1$. In this case, $f_{i} \in Q_{1}(T)$.
c) $\beta_{i} \geq \alpha_{i} \geq 2$. In this case, we have

$$
\begin{aligned}
f_{i} & =\frac{\mathrm{d}^{\beta_{i}-1}}{\mathrm{~d} \xi_{i}^{\beta_{i}-1}}\left(2 \alpha_{i} \xi_{i}\left(\xi_{i}^{2}-1\right)^{\alpha_{i}-1}\right) \\
& =\frac{\mathrm{d}^{\beta_{i}-2}}{\mathrm{~d} \xi_{i}^{\beta_{i}-2}}\left(2 \alpha_{i}\left(\xi_{i}^{2}-1\right)^{\alpha_{i}-1}+4 \alpha_{i}\left(\alpha_{i}-1\right) \xi_{i}^{2}\left(\xi_{i}^{2}-1\right)^{\alpha_{i}-2}\right) \\
& =C_{1} \frac{\mathrm{~d}^{\beta_{i}-2}}{\mathrm{~d} \xi_{i}^{\beta_{i}-2}}\left(\xi_{i}^{2}-1\right)^{\alpha_{i}-1}+C_{2} \frac{\mathrm{~d}^{\beta_{i}-2}}{\mathrm{~d} \xi_{i}^{\beta_{i}-2}}\left(\xi_{i}^{2}-1\right)^{\alpha_{i}-2}
\end{aligned}
$$

where $C_{1}$ and $C_{2}$ are constants. Repeating the same argument, we can read $f_{i}$ as the linear combination of terms satisfying case a) or case b). Then $f_{i} \in G_{i}(T)+Q_{1}(T)$.
d) For $g_{i}$, we have

$$
g_{i}=\frac{1}{2\left(\alpha_{i}+1\right)} \frac{\mathrm{d}^{\beta_{i}+1}}{\mathrm{~d} \xi_{i}^{\beta_{i}+1}}\left(\xi_{i}^{2}-1\right)^{\alpha_{i}+1}
$$

Then $g_{i} \in G_{i}(T)+Q_{1}(T)$ by the discussion from case a) to case c).
Finally, we conclude that $\partial^{\beta} p \in G_{i}(T)+Q_{1}(T)$. Therefore, $\partial^{\beta} p-\Pi_{T, 1} \partial^{\beta} p \in G_{i}(T)+Q_{1}(T)$. Since $\partial^{\beta} p-\Pi_{T, 1} \partial^{\beta} p$ vanishes at the vertices of $T$, we have that $\partial^{\beta} p-\Pi_{T, 1} \partial^{\beta} p \in G_{i}(T)$, and (2.11) is proved.

## 3. $C^{0}$ Nonconforming Plate Elements

In this section, we will focus on the plate bending problem, and consider the nonconforming finite element. Let $n=2$.

Given a triangle $T$, its vertices are denoted by $a_{i}, 1 \leq i \leq 3$. The side of $T$ opposite to $a_{i}$ is denoted by $F_{i}$, its unit outer normal by $\nu_{F_{i}}$ and its measure by $\left|F_{i}\right|, 1 \leq i \leq 3$. Let $\lambda_{1}, \lambda_{2}, \lambda_{3}$ be the barycentric coordinates of $T$. Denote

$$
\left\{\begin{array}{l}
\tilde{q}_{1}=2\left(5\left(\lambda_{1}-\lambda_{1}^{2}-2 \lambda_{2} \lambda_{3}\right)-1\right) \lambda_{1} \lambda_{2} \lambda_{3}  \tag{3.1}\\
\tilde{q}_{2}=2\left(5\left(\lambda_{2}-\lambda_{2}^{2}-2 \lambda_{1} \lambda_{3}\right)-1\right) \lambda_{1} \lambda_{2} \lambda_{3} \\
\tilde{q}_{3}=2\left(5\left(\lambda_{3}-\lambda_{3}^{2}-2 \lambda_{1} \lambda_{2}\right)-1\right) \lambda_{1} \lambda_{2} \lambda_{3}
\end{array}\right.
$$

Set

$$
\begin{equation*}
P_{3}^{+}(T)=P_{3}(T)+\operatorname{span}\left\{\tilde{q}_{1}, \tilde{q}_{2}, \tilde{q}_{3}\right\} \tag{3.2}
\end{equation*}
$$

It is obvious that

$$
\tilde{q}_{1}+\tilde{q}_{2}+\tilde{q}_{3}=-6 \lambda_{1} \lambda_{2} \lambda_{3}
$$

so the dimension of $P_{3}^{+}(T)$ is at most twelve.
The new plate element is defined by $\left(T, P_{T}, D_{T}\right)$ with

1. $T$ is a triangle;
2. $P_{T}=P_{3}^{+}(T)$;
3. the components of $D_{T}(v)$ for any $C^{1}(T)$ are:

$$
\begin{cases}v\left(a_{j}\right), \frac{1}{\left|F_{j}\right|} \int_{F_{j}} \frac{\partial v}{\partial \nu_{F_{j}}} \mathrm{~d} s, & 1 \leq j \leq 3  \tag{3.3}\\ \left(a_{j}-a_{i}\right)^{\mathrm{T}} \nabla v\left(a_{i}\right), & 1 \leq i \neq j \leq 3\end{cases}
$$

where $\nabla$ is the gradient operator.

Define, for $1 \leq i \neq j \neq k \leq 3$,

$$
\left\{\begin{array}{l}
q_{i}=\frac{\tilde{q}_{i}}{\left\|\nabla \lambda_{i}\right\|},  \tag{3.4}\\
p_{i}=3 \lambda_{i}^{2}-2 \lambda_{i}^{3}+\sum_{\substack{1 \leq l \leq 3 \\
l \neq i}} \frac{\nabla \lambda_{i}^{\mathrm{T}} \nabla \lambda_{l}}{\left\|\nabla \lambda_{l}\right\|} q_{l} \\
p_{i j}=\lambda_{i}^{2} \lambda_{j}+10 \lambda_{i}\left(\lambda_{j}-\lambda_{k}\right) \lambda_{1} \lambda_{2} \lambda_{3}
\end{array}\right.
$$

Let $\delta_{i j}$ be the Kronecker delta. It can be verified that $q_{i}, p_{i}, p_{i j} \in P_{3}^{+}(T)$, and

$$
\left\{\begin{array}{lll}
q_{i}\left(a_{k}\right)=0, & \left(a_{l}-a_{k}\right)^{\mathrm{T}} \nabla q_{i}\left(a_{k}\right)=0, & \frac{1}{\left|F_{k}\right|} \int_{F_{k}} \frac{\partial q_{i}}{\partial \nu_{F_{k}}} \mathrm{~d} s=\delta_{i k}  \tag{3.5}\\
p_{i}\left(a_{k}\right)=\delta_{i k}, & \left(a_{l}-a_{k}\right)^{\mathrm{T}} \nabla p_{i}\left(a_{k}\right)=0, & \frac{1}{\left|F_{k}\right|} \int_{F_{k}} \frac{\partial p_{i}}{\partial \nu_{F_{k}}} \mathrm{~d} s=0 \\
p_{i j}\left(a_{k}\right)=0, & \left(a_{l}-a_{k}\right)^{\mathrm{T}} \nabla p_{i j}\left(a_{k}\right)=\delta_{i k} \delta_{j l}, & \frac{1}{\left|F_{k}\right|} \int_{F_{k}} \frac{\partial p_{i j}}{\partial \nu_{F_{k}}} \mathrm{~d} s=0
\end{array}\right.
$$

when $1 \leq i \neq j \leq 3$ and $1 \leq k \neq l \leq 3$. Hence $q_{i}, p_{i}$ and $p_{i j}$ are the nodal basis functions with respect to the degrees of freedom. Therefore, the dimension of $P_{3}^{+}(T)$ is 12 and $D_{T}$ is $P_{T}$-unisolvent.

One can verify that

$$
\begin{equation*}
\int_{F_{k}} \lambda_{l} \frac{\partial q_{i}}{\partial \nu_{F_{k}}} \mathrm{~d} s=\frac{\left|F_{k}\right|}{2} \delta_{i k}, \quad \int_{F_{k}} \lambda_{l} \frac{\partial p_{i}}{\partial \nu_{F_{k}}} \mathrm{~d} s=0 \tag{3.6}
\end{equation*}
$$

when $1 \leq i \leq 3$ and $1 \leq k \neq l \leq 3$, and that

$$
\begin{equation*}
\left.\nabla p_{i j}\right|_{F_{i}} \equiv 0, \quad \int_{F_{k}} \lambda_{i} \frac{\partial p_{i j}}{\partial \nu_{F_{k}}} \mathrm{~d} s=-\int_{F_{k}} \lambda_{l} \frac{\partial p_{i j}}{\partial \nu_{F_{k}}} \mathrm{~d} s=\frac{\left|F_{k}\right|}{12} \nu_{F_{k}}^{\mathrm{T}} \nabla \lambda_{j} \tag{3.7}
\end{equation*}
$$

when $1 \leq i \neq j \leq 3,1 \leq k \neq i \leq 3$ and $1 \leq l \neq k, i \leq 3$.
Given $p \in P_{3}^{+}(T)$, it can be written as

$$
p=\sum_{1 \leq i \leq 3} p\left(a_{i}\right) p_{i}+\sum_{1 \leq i \leq 3} \frac{1}{\left|F_{i}\right|} \int_{F_{i}} \frac{\partial p}{\partial \nu_{F_{i}}} \mathrm{~d} F_{i} q_{i}+\sum_{1 \leq i \neq j \leq 3}\left(a_{j}-a_{i}\right)^{\mathrm{T}} \nabla p\left(a_{i}\right) p_{i j}
$$

Then for $1 \leq i \neq j \neq k \leq 3$, it can be computed by (3.6), (3.7) and the above equality that

$$
\begin{equation*}
\frac{1}{\left|F_{i}\right|} \int_{F_{i}} \lambda_{j} \frac{\partial p}{\partial \nu_{F_{i}}} \mathrm{~d} s=\frac{1}{12}\left(\frac{\partial p}{\partial \nu_{F_{i}}}\left(a_{j}\right)-\frac{\partial p}{\partial \nu_{F_{i}}}\left(a_{k}\right)\right)+\frac{1}{2\left|F_{i}\right|} \int_{F_{i}} \frac{\partial p}{\partial \nu_{F_{i}}} \mathrm{~d} s \tag{3.8}
\end{equation*}
$$

Given any edge $F$ of $T \in \mathcal{T}_{h}$, denote its unit outer normal by $\nu_{F}$. For any $v \in L^{2}(\Omega)$ with $\left.v\right|_{\tilde{T}} \in H^{1}(\tilde{T}), \forall \tilde{T} \in \mathcal{T}_{h}$, we define the jump of $v$ across $F$ as follows:

$$
[v]_{F}=\left.v\right|_{T}-\left.v\right|_{T^{\prime}}
$$

if $F=T \cap T^{\prime}$ for some other $T^{\prime} \in \mathcal{T}_{h}$, and

$$
[v]_{F}=\left.v\right|_{T}
$$

if $F=T \cap \partial \Omega$.
For the new element, define the corresponding finite element spaces $V_{h}$ and $V_{h 0}$ as follows. $V_{h}=\left\{v \in L^{2}(\Omega)|v|_{T} \in P_{3}^{+}(T), \forall T \in T_{h}, v\right.$ and $\nabla v$ are continuous at all vertices of elements in $\mathcal{T}_{h}$, and for any edge $F$ of $T$ with $F \not \subset \partial \Omega$ the integral average of $\nu_{F}^{\mathrm{T}}[\nabla v]_{F}$ over $F$ is zero $\}$; and $V_{h 0}=\left\{v \in V_{h} \mid v\right.$ and $\nabla v$ vanish at all vertices belonging to $\partial \Omega$, and for any edge $F$ of $T$ with $F \subset \partial \Omega$ the integral average of $\frac{\partial}{\partial \nu_{F}} v$ over $F$ is zero $\}$.

We claim that $V_{h} \subset C^{0}(\bar{\Omega})$ and $V_{h 0} \subset C_{0}^{0}(\Omega)$. Let $v_{h} \in V_{h}, F$ be a common edge of $T, T^{\prime} \in \mathcal{T}_{h}$. By the definition, $\left[v_{h}\right]_{F}$ is in $P_{3}(F)$, and it and its directive derivative along $F$ are zero at two endpoints of $F$. Hence $\left[v_{h}\right]_{F} \equiv 0$, that is, $v_{h} \in C^{0}(\bar{\Omega})$. Similarly, we can show that $v_{h} \in C_{0}^{0}(\Omega)$ when $v_{h} \in V_{h 0}$.

By (3.8), the definitions of $V_{h}$ and $V_{h 0}$ and the fact that $V_{h} \subset C^{0}(\bar{\Omega})$ and $V_{h 0} \subset C_{0}^{0}(\Omega)$, we obtain the following lemma.

Lemma 3.1. If $F$ is a common edge of distinct $T, T^{\prime} \in \mathcal{T}_{h}$, then

$$
\begin{equation*}
\int_{F} p\left[\nabla v_{h}\right]_{F} \mathrm{~d} s=0, \quad \forall p \in P_{1}(F), \quad \forall v_{h} \in V_{h} \tag{3.9}
\end{equation*}
$$

If an edge $F$ of $T \in \mathcal{T}_{h}$ is on $\partial \Omega$ then

$$
\begin{equation*}
\int_{F} p\left[\nabla v_{h}\right]_{F} \mathrm{~d} s=0, \quad \forall p \in P_{1}(F), \quad \forall v_{h} \in V_{h 0} \tag{3.10}
\end{equation*}
$$

From the new plate element given above, we can deduce a new Zienkiewicz-type element. For $1 \leq i \leq 3$, define

$$
\begin{equation*}
\phi_{i}(v)=\frac{1}{\left|F_{i}\right|} \int_{F_{i}} \frac{\partial v}{\partial \nu_{F_{i}}} \mathrm{~d} s-\frac{1}{2} \sum_{1 \leq j \leq 3, j \neq i} \frac{\partial v}{\partial \nu_{F_{i}}}\left(a_{j}\right), \quad \forall v \in C^{1}(T) \tag{3.11}
\end{equation*}
$$

Set

$$
\begin{equation*}
P_{T}^{z}=\left\{p \in P_{3}^{+}(T) \mid \phi_{1}(p)=\phi_{2}(p)=\phi_{3}(p)=0\right\} . \tag{3.12}
\end{equation*}
$$

Observing the fact that $P_{2}(T) \subset P_{3}^{+}(T)$ and $\phi_{1}(p)=\phi_{2}(p)=\phi_{3}(p)=0, \forall p \in P_{2}(T)$, we have $P_{2}(T) \subset P_{T}^{z}$.

The new Zienkiewicz-type element is defined by $\left(T, P_{T}, D_{T}\right)$ with

1. $T$ is a triangle;
2. $P_{T}=P_{T}^{z}$;
3. the components of $D_{T}(v)$ for any $v \in C^{1}(T)$ are:

$$
\begin{cases}v\left(a_{j}\right), & 1 \leq j \leq 3,  \tag{3.13}\\ \left(a_{j}-a_{i}\right)^{\mathrm{T}} \nabla v\left(a_{i}\right), & 1 \leq i \neq j \leq 3\end{cases}
$$

It is easy to verify that $D_{T}$ is $P_{T}$-unisolvent.

For the new Zienkiewicz-type element, the corresponding finite element spaces $V_{h}^{z}$ and $V_{h 0}^{z}$ are defined as follows. $V_{h}^{z}=\left\{v \in L^{2}(\Omega)|v|_{T} \in P_{T}^{z}, \forall T \in T_{h}, v\right.$ and $\nabla v$ are continuous at all vertices of elements in $\left.\mathcal{T}_{h}\right\}, V_{h 0}^{z}=\left\{v \in V_{h}^{z} \mid v\right.$ and $\nabla v$ vanish at all vertices belonging to $\left.\partial \Omega\right\}$.

The difference between the new Zienkiewicz-type element here and the two dimensional one proposed in [14] is their shape function spaces. The consistent term of the element here is of order $\mathcal{O}\left(h^{2}\right)$, while the consistent term of the element given in [14] is of order $\mathcal{O}(h)$.

## 4. Convergence Analysis

Let $f \in L^{2}(\Omega)$. We take the following boundary value problem as example to show the convergent result:

$$
\left\{\begin{array}{l}
(-1)^{m} \Delta^{m} u=f, \quad \text { in } \Omega  \tag{4.1}\\
\left.u\right|_{\partial \Omega}=\left.\frac{\partial u}{\partial \nu}\right|_{\partial \Omega}=\cdots=\left.\frac{\partial^{m-1} u}{\partial \nu^{m-1}}\right|_{\partial \Omega}=0
\end{array}\right.
$$

where $\nu=\left(\nu_{1}, \nu_{2}, \cdots, \nu_{n}\right)^{\mathrm{T}}$ is the unit outer normal to $\partial \Omega$ and $\Delta$ is the standard Laplacian operator. Define

$$
\begin{equation*}
a(u, v)=\sum_{1 \leq j_{1}, \cdots, j_{m} \leq n} \int_{\Omega} \frac{\partial^{m} u}{\partial x_{j_{1}} \cdots \partial x_{j_{m}}} \frac{\partial^{m} v}{\partial x_{j_{1}} \cdots \partial x_{j_{m}}} \mathrm{~d} x . \tag{4.2}
\end{equation*}
$$

Then the weak form of problem (4.1) is: find $u \in H_{0}^{m}(\Omega)$ such that

$$
\begin{equation*}
a(u, v)=(f, v), \quad \forall v \in H_{0}^{m}(\Omega) \tag{4.3}
\end{equation*}
$$

For nonnegative integer $s$ and $\mathcal{T}_{h}$, define

$$
H^{s}\left(\mathcal{T}_{h}\right)=\left\{v \in L^{2}(\Omega)|v|_{T} \in H^{s}(T), \quad \forall T \in \mathcal{T}_{h}\right\}
$$

For $v, w \in H^{m}\left(\mathcal{T}_{h}\right)$, define

$$
\begin{equation*}
a_{h}(v, w)=\sum_{T \in \mathcal{T}_{h}} \sum_{1 \leq j_{1}, \cdots, j_{m} \leq n} \int_{T} \frac{\partial^{m} v}{\partial x_{j_{1}} \cdots \partial x_{j_{m}}} \frac{\partial^{m} w}{\partial x_{j_{1}} \cdots \partial x_{j_{m}}} \mathrm{~d} x \tag{4.4}
\end{equation*}
$$

The finite element method for problem (4.3) is: find $u_{h} \in V_{h 0}$ such that

$$
\begin{equation*}
a_{h}\left(u_{h}, v_{h}\right)=\left(f, v_{h}\right), \quad \forall v_{h} \in V_{h 0} . \tag{4.5}
\end{equation*}
$$

We introduce the following mesh dependent norm $\|\cdot\|_{s, h}$ and semi-norm $|\cdot|_{s, h}$ :

$$
\left\{\begin{array}{ll}
\|v\|_{s, h}=\left(\sum_{T \in \mathcal{T}_{h}}\|v\|_{s, T}^{2}\right)^{1 / 2} \\
|v|_{s, h}=\left(\sum_{T \in \mathcal{T}_{h}}|v|_{s, T}^{2}\right)^{1 / 2}
\end{array} \quad \forall v \in H^{s}\left(\mathcal{T}_{h}\right)\right.
$$

For the nonconforming elements, the basic mathematical theory has been established (see [5, $7,10,13,18]$ ). We can use them to give the convergence analysis of our new elements.

For the finite elements given in previous two sections, one can verify the following statements by their constructions, Lemmas 2.1, 2.2 and 3.1:

- They all have the approximability.
- They all have the superapproximation.
- They all have the weak continuity.
- They all pass the patch test.
- They all pass the generalized patch test.

Then by the result in [10] or by the one in [13] we can obtain the following theorems.
Theorem 4.1. Assume that $m, n \geq 1$. Let $V_{h 0}$ be the finite element space corresponding to the rectangular finite element of order $m$, and let $u$ and $u_{h}$ be the solutions of problems (4.3) and (4.5) respectively. Then

$$
\begin{equation*}
\lim _{h \rightarrow 0}\left\|u-u_{h}\right\|_{m, h}=0 \tag{4.6}
\end{equation*}
$$

Theorem 4.2. Assume that $m=n=2$. Let $V_{h 0}$ be the finite element space corresponding to the new plate element or new Zienkiewicz-type element, and let $u$ and $u_{h}$ be the solutions of problems (4.3) and (4.5) respectively. Then

$$
\begin{equation*}
\lim _{h \rightarrow 0}\left\|u-u_{h}\right\|_{2, h}=0 \tag{4.7}
\end{equation*}
$$

By the result in [13], we know that the error of the rectangular finite element of order $m$ and the new Zienkiewicz-type are all order $\mathcal{O}(h)$. For the new plate element, we can obtain the following theorem by Lemma 3.1 and the usual technique dealing with the consistent term.

Theorem 4.3. Assume that $m=n=2$. Let $V_{h 0}$ be the finite element space corresponding to the new plate element, and let $u$ and $u_{h}$ be the solutions of problems (4.3) and (4.5) respectively. Then there exists a constant $C$ independent of $h$ such that

$$
\begin{equation*}
\left\|u-u_{h}\right\|_{2, h} \leq C h^{2}|u|_{4, \Omega} \tag{4.8}
\end{equation*}
$$

when $u \in H^{4}(\Omega)$. In addition,

$$
\begin{equation*}
\left\|u-u_{h}\right\|_{1, \Omega} \leq C h^{3}|u|_{4, \Omega} \tag{4.9}
\end{equation*}
$$

when $\Omega$ is convex.
Remark 4.1. Let $k \geq 1$. The finite element space $V_{h}$ corresponding to the rectangular element of order $k$ is a subspace of $H^{1}(\Omega)$. Hence the element is convergent with order $\mathcal{O}\left(h^{2 k-1}\right)$ by Lemma 2.1 when it is applied to solving the second-oder problems. In general, the rectangular element of order $k$ is a convergent nonconforming element for the $2 m$-th-oder problem when $k \geq m$, which can be shown by Lemmas 2.1 and 2.2 . In this situation, the finite element space $V_{h 0}$ should be defined accordingly.

## 5. Numerical Examples

In this section, we give some numerical results of the new plate element. Now let $m=n=2$, $\Omega=(0,1) \times(0,1)$ and define

$$
\begin{aligned}
& u_{1}(x)=x_{1}^{2}\left(x_{1}-1\right)^{2} x_{2}^{2}\left(x_{2}-1\right)^{2} \\
& u_{2}(x)=\left(\sin \left(\pi x_{1}\right) \sin \left(\pi x_{2}\right)\right)^{2} \\
& u_{3}(x)=e^{x_{1}+x_{2}}
\end{aligned}
$$



Fig. 5.1. The error: $\left|\Pi_{h} u-u_{h}\right|_{2, h}$

For $1 \leq i \leq 3$, set $f_{i}=\Delta^{2} u_{i}$. Then $u_{i}$ is the solution of problem:

$$
\begin{cases}\Delta^{2} u=f_{i}, & \text { in } \Omega  \tag{5.1}\\ u=u_{i}, \frac{\partial u}{\partial \nu}=\frac{\partial u_{i}}{\partial \nu}, & \text { on } \partial \Omega\end{cases}
$$

Problem (5.1) is a homogeneous Dirichlet boundary value problem when $i=1,2$, and a nonhomogeneous one when $i=3$.

For mesh size $h=2^{-1}, 2^{-2}, \cdots, \Omega$ is divided into $h \times h$ squares, and each square is further divided into two triangles by the diagonal with a negative slash.

Let $\Pi_{h}$ be the interpolation operator corresponding to new plate element and $\mathcal{T}_{h}$, and let $u_{h}$ be the finite element solution corresponding to new plate element and triangulation $\mathcal{T}_{h}$. The numerical results of error term $\left|\Pi_{h} u-u_{h}\right|_{2, h}$ are shown in Fig. 5.1 with respect to mesh size $h$. It is seen that the error terms $\left|\Pi_{h} u-u_{h}\right|_{2, h}$ are of $\mathcal{O}\left(h^{2}\right)$ as $h$ approaches 0 . On the other hand, the interpolation error $\left|u-\Pi_{h} u\right|_{2, h}$ is of order $\mathcal{O}\left(h^{2}\right)$ at least. So that $\left|u-u_{h}\right|_{2, h}$ is at least two order of $h$ as well.

Acknowledgments. The work was supported by the National Natural Science Foundation of China (10871011).

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[^0]:    * Received October 23, 2009 / Revised version received May 17, 2010 / Accepted July 20, 2010 /

    Published online November 20, 2010 /

