

Canard Phenomenon and Dynamics for a Slow-Fast Leslie-Gower Prey-Predator Model with Monod-Haldane Function Response

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Abstract The geometrical singular perturbation theory has been successfully applied in studying a large range of mathematical biological models with different time scales. In this paper, we use the geometrical singular perturbation theory to investigate a slow-fast Leslie-Gower prey-predator model with Monod-Haldane function response and get some interesting dynamical phenomena such as singular Hopf bifurcation, canard explosion phenomenon, relaxation oscillation cycle, heteroclinic and homoclinic orbits and the coexistence of canard cycle and relaxation oscillation cycle.

Keywords Leslie-Gower prey-predator model, slow-fast system, canard explosion phenomenon, relaxation oscillation cycle, heteroclinic orbit, homoclinic orbit

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1. Introduction

Predator-prey model is a typical theme in mathematical biology because of its wide application such as biological invasion of foreign species. In this paper, we mainly investigate the modified Leslie-Gower prey-predator model with Monod-Haldane function response [12, 21, 26] as follows

$$\begin{aligned}\frac{du}{dt} &= u(a_1 - b_1u) - \frac{c_1uv}{u^2 + k_1}, \\ \frac{dv}{dt} &= v\left(a_2 - \frac{c_2v}{u^2 + k_2}\right),\end{aligned}\tag{1.1}$$

where u and v separately represent the amount of prey and predators and all parameters are positive and have the following biological meanings: a_1 and a_2 are the natural growth rate of prey and predator which satisfy the assumption that the natural growth rate of prey a_1 is much larger than that of predators a_2 ; b_1 is the intraspecific competition rate of prey; c_1 measures the reduction of prey due

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to predation; c_2 measures the reduction of predator because of the low density of prey; k_1 and k_2 represent the protection provided by the environment of prey and predators. Note that our assumption is reasonable because the lifespan of predators is very long and they may undergo many different generations of prey, such as hares and lynx, squirrels and coyotes.

For simplicity, using the rescaling transformation

$$\begin{aligned} \bar{t} &= a_1 t, & \bar{u} &= \frac{b_1 u}{a_1}, & \bar{v} &= \frac{c_1 b_1^2 v}{a_1^3}, \\ \epsilon &= \frac{a_2}{a_1}, & \bar{c} &= \frac{c_2 a_1}{c_1 a_2}, & \bar{k}_1 &= \frac{k_1 b_1^2}{a_1^2}, & \bar{k}_2 &= \frac{k_2 b_1^2}{a_1^2} \end{aligned}$$

and dropping the bar notation, we rewrite system (1.1) as the following non-dimensional system

$$\begin{aligned} \frac{du}{dt} &= u \left[1 - u - \frac{v}{u^2 + k_1} \right], \\ \frac{dv}{dt} &= \epsilon v \left(1 - \frac{cv}{u^2 + k_2} \right). \end{aligned} \tag{1.2}$$

Furthermore, with transformations

$$\tilde{t} = \int_0^t \frac{1}{(u^2 + k_1)(u^2 + k_2)} ds$$

and $t = \tilde{t}$, system (1.2) can be rewritten as the following topological equivalent form

$$\begin{aligned} \frac{du}{dt} &= u(u^2 + k_2) [(1 - u)(u^2 + k_1) - v] = f(u, v, \eta), \\ \frac{dv}{dt} &= \epsilon v(u^2 + k_1)(u^2 + k_2 - cv) = \epsilon g(u, v, \eta), \end{aligned} \tag{1.3}$$

where $\eta = (c, k_1, k_2)$ and $0 < \epsilon \ll 1$ under our assumption $a_2 \ll a_1$

It is clear that system (1.3) is a slow-fast system with one slow state variable v and one fast state variable u . Using the change of time scale $\tau = \epsilon t$, the system (1.3) can be transformed as

$$\begin{aligned} \epsilon \frac{du}{d\tau} &= u(u^2 + k_2) [(1 - u)(u^2 + k_1) - v] = f(u, v, \eta), \\ \frac{dv}{d\tau} &= v(u^2 + k_1)(u^2 + k_2 - cv) = g(u, v, \eta). \end{aligned} \tag{1.4}$$

Compared with t , the time scale τ is the slow time scale. Hence, systems (1.3) and (1.4) are separately called the fast system and the slow system and their dynamics are equivalent if $0 < \epsilon \ll 1$.

Letting $\epsilon \rightarrow 0$ in systems (1.3) and (1.4), we get the degenerate system

$$\begin{aligned} 0 &= f(u, v, \eta), \\ \frac{dv}{d\tau} &= g(u, v, \eta), \end{aligned}$$