

## A High Order Central DG method of the Two-Layer Shallow Water Equations

Yongping Cheng<sup>1</sup>, Haiyun Dong<sup>1</sup>, Maojun Li<sup>2,\*</sup> and Weizhi Xian<sup>1</sup>

<sup>1</sup> College of Mathematics and Statistics, Chongqing University, Chongqing, 401331, P.R. China.

<sup>2</sup> School of Mathematical Sciences, University of Electronic Science and Technology of China, Sichuan, 611731, P.R. China.

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**Abstract.** In this paper, we focus on the numerical simulation of the two-layer shallow water equations over variable bottom topography. Although the existing numerical schemes for the single-layer shallow water equations can be extended to two-layer shallow water equations, it is not a trivial work due to the complexity of the equations. To achieve the well-balanced property of the numerical scheme easily, the two-layer shallow water equations are reformulated into a new form by introducing two auxiliary variables. Since the new equations are only conditionally hyperbolic and their eigenstructure cannot be easily obtained, we consider the utilization of the central discontinuous Galerkin method which is free of Riemann solvers. By choosing the values of the auxiliary variables suitably, we can prove that the scheme can exactly preserve the still-water solution, and thus it is a truly well-balanced scheme. To ensure the non-negativity of the water depth, a positivity-preserving limiter and a special approximation to the bottom topography are employed. The accuracy and validity of the numerical method will be illustrated through some numerical tests.

**AMS subject classifications:** 65M60, 65M12

**Key words:** Two-layer shallow water equations, Central DG method, Positivity-preserving and well-balanced, Still-water solution.

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## 1 Introduction

The shallow water equations [19, 20] are widely adopted to model free-surface flows in rivers, channels, flood plains and coastal regions, however, the single-layer shallow water equations have the drawback of missing some physical dynamics in the vertical motion when studying stratified flow motions. Therefore, during the last decades,

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\*Corresponding author. *Email addresses:* cyp@cqu.edu.cn (Y. Cheng), dhy@cqu.edu.cn (H. Dong), limj@cqu.edu.cn (M. Li), wasxxwz@163.com (W. Xian)

multi-layer shallow water equations have been attracted more attention and have become an important tool to study stratified flow motions, such as salinity-driven exchange flow motions [12], wave-mud interactions [17] and chaotic mixing of particles in layered flows [24].

In this paper, we focus on the numerical simulation of the two-layer shallow water equations, which are given by

$$\left\{ \begin{aligned} \frac{\partial h_1}{\partial t} + \frac{\partial(h_1 u_1)}{\partial x} &= 0, \\ \frac{\partial(h_1 u_1)}{\partial t} + \frac{\partial(h_1 u_1^2 + \frac{1}{2} g h_1^2)}{\partial x} &= -g h_1 \frac{\partial(h_2 + b)}{\partial x}, \\ \frac{\partial h_2}{\partial t} + \frac{\partial(h_2 u_2)}{\partial x} &= 0, \\ \frac{\partial(h_2 u_2)}{\partial t} + \frac{\partial(h_2 u_2^2 + \frac{1}{2} g h_2^2)}{\partial x} &= -g h_2 \frac{\partial b}{\partial x} - g h_2 \frac{\rho_1}{\rho_2} \frac{\partial h_1}{\partial x}, \end{aligned} \right. \tag{1.1}$$

in one-dimensional (1-D) space, and

$$\left\{ \begin{aligned} \frac{\partial h_1}{\partial t} + \frac{\partial(h_1 u_1)}{\partial x} + \frac{\partial(h_1 v_1)}{\partial y} &= 0, \\ \frac{\partial(h_1 u_1)}{\partial t} + \frac{\partial(h_1 u_1^2 + \frac{1}{2} g h_1^2)}{\partial x} + \frac{\partial(h_1 u_1 v_1)}{\partial y} &= -g h_1 \frac{\partial(h_2 + b)}{\partial x}, \\ \frac{\partial(h_1 v_1)}{\partial t} + \frac{\partial(h_1 u_1 v_1)}{\partial x} + \frac{\partial(h_1 v_1^2 + \frac{1}{2} g h_1^2)}{\partial y} &= -g h_1 \frac{\partial(h_2 + b)}{\partial y}, \\ \frac{\partial h_2}{\partial t} + \frac{\partial(h_2 u_2)}{\partial x} + \frac{\partial(h_2 v_2)}{\partial y} &= 0, \\ \frac{\partial(h_2 u_2)}{\partial t} + \frac{\partial(h_2 u_2^2 + \frac{1}{2} g h_2^2)}{\partial x} + \frac{\partial(h_2 u_2 v_2)}{\partial y} &= -g h_2 \frac{\partial b}{\partial x} - g h_2 \frac{\rho_1}{\rho_2} \frac{\partial h_1}{\partial x}, \\ \frac{\partial(h_2 v_2)}{\partial t} + \frac{\partial(h_2 u_2 v_2)}{\partial x} + \frac{\partial(h_2 v_2^2 + \frac{1}{2} g h_2^2)}{\partial y} &= -g h_2 \frac{\partial b}{\partial y} - g h_2 \frac{\rho_1}{\rho_2} \frac{\partial h_1}{\partial y}, \end{aligned} \right. \tag{1.2}$$

in two-dimensional (2-D) space. Herein, the subscripts 1 and 2 denote the upper and the lower layer in the system respectively (see Fig. 1). In these equations,  $\rho_i$  is a constant representing the density of the  $i$ th layer with  $0 < \rho_1 < \rho_2$ ,  $h_i$  denotes the height of the  $i$ th layer,  $u_i, v_i$  are the vertically averaged horizontal velocity of the  $i$ th layer in the  $x$ - and  $y$ -directions, respectively,  $i = 1, 2$ ,  $b$  is the bottom topography and  $g$  represents the gravitational constant.

Similar to single-layer shallow water equations, these equations admit still-water solutions

$$u_1 = u_2 = 0, \quad h_1 + h_2 + b = \text{constant}, \quad h_2 + b = \text{constant} \tag{1.3}$$

for 1-D case and

$$u_1 = u_2 = v_1 = v_2 = 0, \quad h_1 + h_2 + b = \text{constant}, \quad h_2 + b = \text{constant} \tag{1.4}$$