

Quasi-periodic Solutions of the General Nonlinear Beam Equations*

GAO YI-XIAN^{1,2}

(1. *College of Mathematics and Statistics, Northeast Normal University,
Changchun, 130024*)

(2. *Key Laboratory of Symbolic Computation and Knowledge Engineering of
Ministry of Education, Jilin University, Changchun, 130012*)

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Abstract: In this paper, one-dimensional (1D) nonlinear beam equations of the form

$$u_{tt} - u_{xx} + u_{xxxx} + mu = f(u)$$

with Dirichlet boundary conditions are considered, where the nonlinearity f is an analytic, odd function and $f(u) = O(u^3)$. It is proved that for all $m \in (0, M^*] \subset \mathbf{R}$ (M^* is a fixed large number), but a set of small Lebesgue measure, the above equations admit small-amplitude quasi-periodic solutions corresponding to finite dimensional invariant tori for an associated infinite dimensional dynamical system. The proof is based on an infinite dimensional KAM theory and a partial Birkhoff normal form technique.

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1 Introduction and Main Result

Consider the general nonlinear beam equations of the form

$$u_{tt} - u_{xx} + u_{xxxx} + mu = f(u) \tag{1.1}$$

on the finite x -interval $[0, \pi]$ with Dirichlet boundary conditions

$$u(t, 0) = u(t, \pi) = u_{xx}(t, 0) = u_{xx}(t, \pi) = 0, \tag{1.2}$$

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where the parameter $m \in (0, M^*] \subset \mathbf{R}$, the nonlinearity f is assumed to be real analytic in u and of the form

$$f(u) = au^3 + \sum_{n \geq 5} f_n u^n, \quad a \neq 0. \quad (1.3)$$

We study the equations of the form (1.1) as a Hamiltonian system on

$$\mathcal{P} = H_0^1([0, \pi]) \times L^2([0, \pi])$$

with coordinates u and $v = u_t$. Then the Hamiltonian is

$$H = \frac{1}{2} \langle v, v \rangle + \frac{1}{2} \langle Au, u \rangle + \int_0^\pi g(u) dx, \quad (1.4)$$

where

$$A = \frac{d^4}{dx^4} - \frac{d^2}{dx^2} + m, \quad g = \int_0^\pi -f(s) ds, \quad (1.5)$$

and $\langle \cdot, \cdot \rangle$ denotes the usual scalar product in L^2 . Then (1.1) can be written in the form

$$u_t = \frac{\partial H}{\partial v} = v, \quad v_t = -\frac{\partial H}{\partial u} = -Au - f(u). \quad (1.6)$$

Let

$$\phi_j(x) = \sqrt{\frac{2}{\pi}} \sin jx, \quad \lambda_j = \sqrt{j^4 + j^2 + m}, \quad j = 1, 2, \dots$$

be the basic modes and frequencies of the linear equation

$$u_{tt} - u_{xx} + u_{xxxx} + mu = 0$$

with Dirichlet boundary conditions (1.2). Then every solution of the linear equation is the superposition of their harmonic oscillations and of the form

$$u(t, x) = \sum_{j \geq 1} q_j(t) \phi_j(x), \quad q_j(t) = \sqrt{I_j} \cos(\lambda_j t + \theta_j),$$

with amplitudes $I_j \geq 0$ and initial phases θ_j . The motions are periodic or quasi-periodic, respectively, depending on whether one or finitely many eigenfunctions are excited. In particular, for every choice

$$J = \{j_1 < j_2 < \dots < j_n\} \subset \mathbf{N}$$

of finitely many modes there exists an invariant $2n$ -dimensional linear subspace E_J which is completely foliated into rotational tori with frequencies $\lambda_{j_1}, \dots, \lambda_{j_n}$:

$$E_J = \{(u, v) = (q_1 \phi_{j_1} + \dots + q_n \phi_{j_n}, p_1 \phi_{j_1} + \dots + p_n \phi_{j_n})\} = \bigcup_{I \in P^n} \mathcal{T}_J(I),$$

where

$$P^n = \{I \in \mathbf{R}^n : I_j > 0, 1 \leq j \leq n\}$$

is the positive quadrant in \mathbf{R}^n and

$$\mathcal{T}_J(I) = \{(u, v) : q_j^2 + \lambda_j^{-2} p_j^2 = I_j, 1 \leq j \leq n\},$$

by using the above representations of u and v . In addition, such a torus is linearly stable, and all solutions have zero Lyapunov exponents.

Upon restoration of the nonlinearity f , we show that there exist a Cantor set $\mathcal{O} \subset P^n$, a family of n -tori

$$\mathcal{T}_J[\mathcal{O}] = \bigcup_{I \in \mathcal{O}} \mathcal{T}_J(I) \subset E_J \quad \text{over } \mathcal{O},$$