

# Regularity for 3-D MHD Equations in Lorentz Space $L^{3,\infty}$

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**Abstract.** The regularity for 3-D MHD equations is considered in this paper, it is proved that the solutions  $(v, B, p)$  are Hölder continuous if the velocity field  $v \in L^\infty(0, T; L_x^{3,\infty}(\mathbb{R}^3))$  with local small condition

$$r^{-3} \left| \left\{ x \in B_r(x_0) : |v(x, t_0)| > \varepsilon r^{-1} \right\} \right| \leq \varepsilon$$

and the magnetic field  $B \in L^\infty(0, T; VMO^{-1}(\mathbb{R}^3))$ .

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**Key words:** Lorentz space, backward uniqueness, MHD equations.

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## 1 Introduction

In this paper, we will consider the following MHD equations in  $\mathbb{R}^3 \times (0, T)$ :

$$\begin{cases} \partial_t v + v \cdot \nabla v - \Delta v + \nabla p = B \cdot \nabla B, \\ \partial_t B - \Delta B + v \cdot \nabla B - B \cdot \nabla v = 0, \\ \operatorname{div} v = 0, \quad \operatorname{div} B = 0, \end{cases} \quad (1.1)$$

where  $v$  is the fluid velocity field and  $B$  is the magnetic field. The function  $p$  describes the scalar pressure. The MHD equations (1.1) reduce to the incompressible Navier-Stokes equations when there is no electromagnetic field, that is,  $B \equiv 0$ .

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There is a very rich literatures dedicated to the mathematical study of the Navier-Stokes system by many researchers; see, for example, [1,3–5,8–10,13–15,18,20,22]. Compared with Navier-Stokes equation, the MHD equation (1.1) is a combination of the incompressible Navier-Stokes equations of fluid dynamics and Maxwell's equations of electromagnetism. Such a system depicts the motion of many conducting incompressible immiscible fluids without surface tension under the action of a magnetic field.

For a point  $z = (x, t) \in \mathbb{R}^3 \times \mathbb{R}_+$ , we denote

$$B_r(x) := \{y \in \mathbb{R}^3 : |y - x| < r\}, \quad Q_r(z) := B_r(x) \times (t - r^2, t), \quad Q_r = Q_r(0, 0),$$

and define the space

$$L^{q,p}(Q_r) = L^q(t - r^2, t; L^p(B_r(x)))$$

with the norm

$$\|u\|_{L^{p,q}(Q_r(z))}^q = \int_{t-r^2}^t \|u(t)\|_{L^p(B_r(x))}^q dt.$$

Let  $\Omega \subset \mathbb{R}^3$ . The Lorentz space  $L^{(p,q)}(\Omega)$  with  $p, q \in (0, \infty)$  is the set of measurable functions  $f$  on  $\Omega$  such that the following quasi-norm is bounded:

$$\|f\|_{L^{(p,q)}(\Omega)} := \begin{cases} \left( p \int_0^\infty \alpha^q |\{x \in \Omega : |f(x)| > \alpha\}|^{\frac{q}{p}} \frac{d\alpha}{\alpha} \right)^{\frac{1}{q}}, & \text{if } q < \infty, \\ \sup_{\alpha > 0} \alpha |\{x \in \Omega : |f(x)| > \alpha\}|^{\frac{1}{p}}, & \text{if } q = \infty. \end{cases}$$

Lorentz spaces can be used to capture logarithmic singularities. For example, in  $\mathbb{R}^3$ , for any  $\beta > 0$  we have

$$|x|^{-1} \left( \frac{\log|x|}{2} \right)^{-\beta} \in L^{(3,q)}(\mathbb{R}^3), \quad \text{if and if } q > \frac{1}{\beta}.$$

It is known that  $L^{(p,p)}(\Omega) = L^p(\Omega)$  and  $L^{(p,q_1)}(\Omega) \subset L^{(p,q_2)}(\Omega)$  while  $q_1 \leq q_2$ . If  $|\Omega|$  is finite then  $L^{(p,q)}(\Omega) \subset L^r(\Omega)$  for all  $0 < q \leq \infty$  and  $0 < r < p$ ,

$$\|g\|_{L^r(\Omega)} \leq |\Omega|^{\frac{1}{r} - \frac{1}{p}} \|g\|_{L^{(p,q)}(\Omega)}. \quad (1.2)$$

In the classical literatures, Leray [9] and Hopf [8] established the existence weak solution to the Navier Stokes system. However, the smoothness of the Leray-Hopf weak solution is still a challenging open problem. Serrin [18] proved