

A SPLITTING ITERATION METHOD FOR DOUBLE X_0 -BREAKING BIFURCATION POINTS*

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Abstract

A splitting iteration method is proposed to compute double X_0 -breaking bifurcation points. The method will reduce the computational work and storage, it converges linearly with an adjustable speed. Numerical computation shows the effectiveness of splitting iteration method.

Key words: Double X_0 -breaking bifurcation point, splitting iteration method, extended system

1. Introduction

Consider the following two-parameter dependent nonlinear problem

$$f(x, \lambda, \mu) = 0, \quad f : X \times R^2 \rightarrow X, \quad (1.1)$$

where $X = R^n$, λ, μ are real parameters, $f \in C^r (r \geq 3)$, $D_x f_0 (\equiv D_x f(x_0, \lambda_0, \mu_0))$ is a Fredholm map with index zero. One of our main assumptions, which arise in many applications^[1,2,5-7], is that f satisfies Z_2 -symmetry: there exists a linear operator $S : X \rightarrow X$ such that

$$S \neq I, S^2 = I, S f(x, \lambda, \mu) = f(Sx, \lambda, \mu), \quad \forall (x, \lambda, \mu) \in X \times R^2 \quad (1.2)$$

It is well-known that X has the following natural decomposition:

$$X = X_s \oplus X_a,$$

where

$$X_s = \{x \in X : Sx = x\}, \quad X_a = \{x \in X : Sx = -x\}$$

are the set of symmetric elements and the set of anti-symmetric elements respectively^[7].

We also assume that there is an invariant subspace $X_0 \subset X_s$ such that

$$f(x, \lambda, \mu) \in X_0, \quad \forall (x, \lambda, \mu) \in X_0 \times R^2. \quad (1.3)$$

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The usual and best case is $X_0 = \{0\}$. We call (x_0, λ_0, μ_0) a singular point of (1.1) if $f(x_0, \lambda_0, \mu_0) = 0$ and $\dim(\text{Null}(f_x(x_0, \lambda_0, \mu_0))) \geq 1$. In this paper, we are concerned with double X_0 -breaking bifurcation points (x_0, λ_0, μ_0) in the sense that^[6]

$$f(x_0, \lambda_0, \mu_0) = 0, x_0 \in X_0, \quad (1.4a)$$

$$\text{Null}(D_x f_0) = \text{span} \{\phi_s, \phi_a\}, \phi_s \in X_s, \phi_s \notin X_0, \phi_a \in X_a, \quad (1.4b)$$

$$\text{Range}(D_x f_0) = \{y \in X : \langle \psi_s, y \rangle = \langle \psi_a, y \rangle = 0\}, \psi_s \in X_s, \psi_a \in X_a, \quad (1.4c)$$

$$\langle \psi_s, D_\lambda f_0 \rangle = \langle \psi_s, D_\mu f_0 \rangle = 0. \quad (1.4d)$$

In addition, as is common, we assume that $\langle \psi_r, \phi_r \rangle = \langle \psi_r, \psi_r \rangle = \langle \phi_r, \phi_r \rangle = 1$, $r = s, a$, $\langle \psi_r, \phi_\delta \rangle = 0$, $(r, \delta) = (s, a)$ or (a, s) . X_0 -breaking bifurcation point is one of the three most important kinds of bifurcation points (the others are turning points and pitchfork points^[2]). For the computation of double X_0 -breaking bifurcation points of (1.1), Werner^[6] proposed a regular extended system which is a direct method and is at least three times larger than the original equation (1.1). Here we will propose a splitting iteration method. The method produces smaller systems and it could simultaneously compute the point (x_0, λ_0, μ_0) , the null vectors of $D_x f_0, D_x f_0^*$ in a coupled way. This method converges linearly with an adjustable speed and its computational cost at each iteration step remains the same level as that for the regular solution of (1.1)^[3,4].

We will construct small extended systems in section 2, then propose the splitting iteration method in section 3. Numerical examples are given in section 4 to show the effectiveness of the method.

2. Extended Systems

First, we introduce the following lemma, which could be proved directly by differentiating (1.2).

Lemma 1. $\forall x \in X_0, \lambda \in R, \mu \in R,$

(i) $f(x, \lambda, \mu), D_\lambda f(x, \lambda, \mu), D_\mu f(x, \lambda, \mu) \in X_0;$

(ii) X_0, X_s and X_a are invariant subspaces of $D_x f(x, \lambda, \mu), D_{x\lambda} f(x, \lambda, \mu), D_{x\mu} f(x, \lambda, \mu);$

(iii) $\forall v, w \in X_0, D_{xx} f(x, \lambda, \mu)vw \in X_0;$

(iv) $\forall v \in X_s, w \in X_s,$ or $\forall v \in X_a, w \in X_a, D_{xx} f(x, \lambda, \mu)vw \in X_s;$

(v) $\forall v \in X_s, w \in X_a, D_{xx} f(x, \lambda, \mu)vw \in X_a.$

It follows from Lemma 1 and (1.4d) that there exist $v_0, u_0 \in X_0$ such that

$$D_x f_0 v_0 + D_\lambda f_0 = 0, \quad (2.1a)$$

$$D_x f_0 u_0 + D_\mu f_0 = 0, \quad (2.1b)$$

and hence we could introduce the following notations

$$A_r := \langle \psi_r, (D_{xx} f_0 v_0 + D_{x\lambda} f_0) \phi_r \rangle, \quad (2.2a)$$

$$B_r := \langle \psi_r, (D_{xx} f_0 u_0 + D_{x\mu} f_0) \phi_r \rangle, r = s, a. \quad (2.2b)$$